

Extention of Newton's Dynamical Field Equation in the Gravitational Field of Spherical Mass

Lucas W. L.^{1,*}, Joseph I. K.², Pius I. P.³, Ewa I. I.¹, Sarki M. U.¹

¹Department of Physics, Nasarawa State University, Keffi, Keffi, Nigeria

²Department of Mathematics, Nasarawa State University, Keffi, Keffi, Nigeria

³University of Glasgow, School of Physics and Astronomy, Scotland UK

Abstract Despite the impressive success of the Newtonian dynamical theory of gravitation in explaining the observed gravitational phenomena, the theory reminds logically incomplete. In this article, the golden Riemannian geometry of space-time was applied to construct an extended Newton's dynamical field equation and the gravitational scalar potential exterior and interior to the gravitational field of static homogeneous spherical mass. The obtained results reduce in the limit c^0 , to the corresponding pure Newtonian expressions. Results obtained contain post-Newton correction terms to the order of c^{-2} . The consequences of the additional correctional terms are that they can be applied to predict the existence of gravitational waves and their propagation in space and to derive the anomalous orbital precession of the orbit of the planet as well as the gravitational shift by the Sun as excellently as Einstein's geometrical theory of gravitation. Also, it can be applied to examine the theoretical analysis of the motion of particles of non-zero rest masses and in a simple linear harmonic oscillator to determine its energy levels.

Keywords Newton's Theory, Riemannian geometry, Scalar Potential, Precession, Red Shift, Harmonic Oscillator, and Energy Levels

1. Introduction

Gravity is known to be the least understood of all the fundamental forces in nature. Interestingly, mass and space which are governed by gravitation are the basic building blocks of the universe [1,5].

Around the year 300 BC, the Greek Philosopher and Mathematician called Euclid introduced his theory of geometry for space-time without the influence and interaction of gravitational fields which is popularly known today as Euclidian geometry [2,9,1,5,10].

In the year 1854, the German Mathematician called George Friedrich Bernard Riemann introduced his principle of geometry for space-time with the influence and interaction of gravitational fields which is known today as the Riemannian geometry. This theory was formulated in terms of the mathematical quantity called the metric tensor of space-time with the influence and interaction of gravitational fields. It is important to note that George Friedrich Bernard Riemann never specified any explicit expression for the metric tensor or tensors for any given gravitational field [2,3,4,6].

In the year 1686, Isaac Newton published his dynamical

theory of gravitation based on Euclidean geometry. According to this theory, all interactions through nature manifest through force. This theory was successful in explaining the gravitational phenomena on earth and the experimental facts of the solar system. It is a well-known fact that Newton's dynamical theory of gravitation (NDTG) could not explain the anomalous orbital precession of the orbit of the planet as well as the gravitational redshift by the Sun.

In the year 1905, Einstein's published his geometrical theory of gravitation (EGTG) which is popularly known today as General Relativity (GR). According to this theory, gravitation is not due to force but due to the geometrical curving of space and time. This theory offered the resolution of the anomalous orbital precession of the orbit of the planet as well as the gravitational redshift by the Sun. Despite the famous test of General Relativity (GR), several Physicists have continued to hold on to the view that Newton's dynamical theory of gravitation can be extended in such a way as to account satisfactorily for the experimental data. Thus far, the General Theory of Relativity has been the most fundamental and successful theory of Physics that successfully explains the nature of gravity. If we better understand the nature of mass and space, we may be able to do things previously undreamed of, so far studies of General Relativity and Gravity have yielded atomic clocks, guidance systems for spacecraft, and the Global Positioning Systems

* Corresponding author:

lumbilucasw@nsuk.edu.ng (Lucas W. L.)

Received: Oct. 18, 2021; Accepted: Mar. 6, 2022; Published: Jun. 13, 2022

Published online at <http://journal.sapub.org/ijtmp>

(GPS). We cannot foresee all that can come from a better understanding of space-time and mass-energy, but a theory about these fundamental subjects must be thoroughly examined if we are to use it to our advantage [2,9,1,5,10].

In the year 2010, Howusu, S. X. K. developed a hybrid theory of gravitation that combines Newton's dynamical theory and Einstein's geometrical theory into a Riemannian theory. He postulated that Newton's dynamical theory of gravitation can be extended using the Great, Golden, and Super Golden Riemannian Laplacian operators [3,4].

It is interesting and instructive to note that Howusu never showed how the Riemannian Laplacian operators can be applied to extend Newton's dynamical theory of gravitation to account satisfactorily with experimental results and data.

This research will substantially extend Newton's dynamical theory of gravitation using a unique approach of golden Riemannian geometry of space-time to obtain an extended Newtonian dynamical field equation and the gravitational scalar potential exterior and interior in the gravitational field of static homogeneous spherical mass.

2. Theoretical Analysis

Consider a static homogeneous spherical body of radius R and total mass M . The golden Riemannian metric tensor in the gravitational field exterior to a spherical mass is given explicitly by [3,4]

$$g_{00}(r, \theta, \phi, x^0) = -\left\{1 - \frac{2}{c^2}f(r, \theta, \phi, x^0)\right\} \quad (1)$$

$$g_{11}(r, \theta, \phi, x^0) = \left\{1 - \frac{2}{c^2}f(r, \theta, \phi, x^0)\right\}^{-1} \quad (2)$$

$$g_{22}(r, \theta, \phi, x^0) = r^2 \left\{1 - \frac{2}{c^2}f(r, \theta, \phi, x^0)\right\}^{-1} \quad (3)$$

$$g_{33}(r, \theta, \phi, x^0) = r^2 \sin^2 \theta \left\{1 - \frac{2}{c^2}f(r, \theta, \phi, x^0)\right\}^{-1} \quad (4)$$

$$g_{\mu\nu} = 0; \text{ otherwise} \quad (5)$$

where, f is an extended Riemannian gravitational scalar potential exterior to the spherical body and (ct, r, θ, ϕ) is the space-time coordinate with $(x^0 = ct)$.

The well-known Riemannian Laplacian operator $\nabla_R^2 f$ is given explicitly as [11,8]

$$\nabla_R^2 f = \frac{1}{\sqrt{g}} \frac{\partial}{\partial x^\alpha} \left(\sqrt{g} g^{\alpha\beta} \frac{\partial}{\partial x^\beta} \right) \quad (6)$$

or more explicitly

$$\begin{aligned} \nabla_R^2 f = & \frac{1}{\sqrt{g}} \frac{\partial}{\partial x^1} \left(\sqrt{g} g^{11} \frac{\partial}{\partial x^1} \right) + \frac{1}{\sqrt{g}} \frac{\partial}{\partial x^2} \left(\sqrt{g} g^{22} \frac{\partial}{\partial x^2} \right) + \\ & \frac{1}{\sqrt{g}} \frac{\partial}{\partial x^3} \left(\sqrt{g} g^{33} \frac{\partial}{\partial x^3} \right) + \frac{1}{\sqrt{g}} \frac{\partial}{\partial x^0} \left(\sqrt{g} g^{00} \frac{\partial}{\partial x^0} \right) \end{aligned} \quad (7)$$

The covariant metric tensor for spherical polar coordinate is given explicitly by

$$g = \begin{vmatrix} g_{11} & 0 & 0 & 0 \\ 0 & g_{22} & 0 & 0 \\ 0 & 0 & g_{33} & 0 \\ 0 & 0 & 0 & g_{00} \end{vmatrix} \quad (8)$$

Substituting equations (1) - (5) into (8), we get

$$g = -r^4 \sin^2 \theta \left\{1 - \frac{2}{c^2}f(r, \theta, \phi, x^0)\right\}^{-2} \quad (9)$$

and

$$\sqrt{g} = -r^2 \sin \theta \left\{1 - \frac{2}{c^2}f(r, \theta, \phi, x^0)\right\}^{-1} \quad (10)$$

Similarly, the contravariant metric tensor is given as

$$g^{\alpha\beta} = \frac{1}{g_{\alpha\beta}} \quad (11)$$

where $\alpha, \beta = 0, 1, 2, 3$.

Hence, substituting (1) - (5) into (8), we obtain

$$g^{00} = -\left\{1 - \frac{2}{c^2}f(r, \theta, \phi, x^0)\right\}^{-1} \quad (12)$$

$$g^{11} = \left\{1 - \frac{2}{c^2}f(r, \theta, \phi, x^0)\right\} \quad (13)$$

$$g^{22} = \frac{1}{r^2} \left\{1 - \frac{2}{c^2}f(r, \theta, \phi, x^0)\right\} \quad (14)$$

$$g^{33} = \frac{1}{r^2 \sin^2 \theta} \left\{1 - \frac{2}{c^2}f(r, \theta, \phi, x^0)\right\} \quad (15)$$

Substituting equation (10) and equations (12) - (15) into (7) we obtain

$$\begin{aligned} \nabla_R^2 f = & \frac{1}{r^2} \left(1 + \frac{2f}{c^2}\right) \left\{ \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \right) \right\} f^+(r) + \frac{1}{r^2 \sin \theta} \\ & \left(1 + \frac{2f}{c^2}\right) \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial}{\partial \theta} \right) + \frac{1}{r^2 \sin \theta} \left(1 + \frac{2f}{c^2}\right) \frac{\partial}{\partial \phi} \left(\frac{1}{\sin \theta} \frac{\partial}{\partial \phi} \right) \\ & + \frac{1}{r^2 \sin \theta} \left(1 + \frac{2f}{c^2}\right) \frac{\partial}{\partial x^0} \left(r^2 \sin \theta \left(1 + \frac{2f}{c^2}\right)^{-2} \frac{\partial}{\partial x^0} \right) = \\ & \begin{cases} 0; & r > R \\ 4\pi G \rho_0; & r < R \end{cases} \end{aligned} \quad (16)$$

where

$x^0 = ct$ is the space time coordinate

$\nabla_R^2 f$ is the Riemannian Laplacian operator

$f(x^0, r, \theta, \phi)$ is the complete generalized gravitational scalar potential exterior to the body.

By the symmetry of the distribution of mass, it follows that the gravitational field will depend only on the radial coordinate r .

Therefore,

$$\frac{\partial f}{\partial \theta} = 0 \quad (17)$$

$$\frac{\partial f}{\partial \phi} = 0 \quad (18)$$

$$\frac{\partial f}{\partial x^0} = 0 \quad (19)$$

Consequently, equation (16) reduces

$$\nabla_R^2 f = \frac{1}{r^2} \left(1 + \frac{2f}{c^2}\right) \left\{ \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \right) \right\} f^+(r) = \begin{cases} 0; & r > R \\ 4\pi G \rho_0; & r < R \end{cases} \quad (20)$$

or more explicitly

$$f'' + \frac{2}{r} f' + \frac{2}{c^2} f f'' + \frac{4}{c^2 r} f f' = \begin{cases} 0; & r > R \\ 4\pi G \rho_0; & r < R \end{cases} \quad (21)$$

Equation (21) is an extended Newtonian dynamical gravitational field equation for spherical fields based on the golden Riemannian geometry of space-time.

This equation:

- Reduces in the limit of c^0 , to the corresponding pure

Newton's dynamical gravitational field equation and in the order of c^{-2} , it contains unknown post-Newton correction terms which are uncovered for theoretical development and experimental investigations and applications.

- Contain $\left(1 + \frac{2}{c^2}f\right)$ which are not found in the existing well-known Newtonian dynamical or Einstein geometrical field equations.
- Contain the time - component which is not found in the existing well - known Newton's dynamical field equation.

The immediate consequences of the additional correction terms obtained in the extended Newtonian field equation are that it will predict:

- Additional correction terms to the gravitational field of all massive bodies and by extension to the gravitational scalar potential exterior and interior.
- Predict the existence of gravitational waves and their propagation in space with the phase velocity, v_p given by $v_p = c$ where, c is the speed of light in vacuum.

3. Construction of an Extended Newtonian Gravitational Scalar Potential Exterior and Interior for Static Homogeneous Spherical Mass

Let us seek the solution for the exterior field equation (21) as

$$f^+(r) = \frac{A_1}{r} + \frac{A_2}{r^2} + \frac{A_3}{r^3} + \frac{A_4}{r^4} + \dots \quad (22)$$

$$f'(r) = -\frac{A_1}{r^2} - \frac{2A_2}{r^3} - \frac{3A_3}{r^4} - \frac{4A_4}{r^5} + \dots \quad (23)$$

$$f''(r) = \frac{2A_1}{r^3} + \frac{6A_2}{r^4} + \frac{12A_3}{r^5} + \frac{12A_4}{r^6} + \dots \quad (24)$$

$$ff'' = \frac{2A_1^2}{r^4} + \frac{8A_1A_2}{r^5} + \frac{14A_1A_3}{r^6} + \frac{22A_1A_4}{r^7} + \frac{6A_2^2}{r^6} + \frac{18A_2A_3}{r^7} + \dots \quad (25)$$

$$ff' = -\frac{A_1^2}{r^4} - \frac{3A_1A_2}{r^5} - \frac{4A_1A_3}{r^6} - \frac{2A_2^2}{r^6} - \frac{4A_2A_3}{r^7} - \frac{3A_3^2}{r^8} + \dots \quad (26)$$

Substituting equations (22) - (26) into equation (21) and equating coefficients of r^{-4} , r^{-5} , r^{-6} and r^{-7} on both sides gives

$$A_1 = A_2 = A_3 = 0 = \text{Arbitrary constant} \quad (27)$$

$$A_4 = -\frac{4A_1^2}{3c^2} \quad (28)$$

Hence, the general solution for the exterior field equation (21) is given by

$$f^+(r) = \frac{A_1}{r} + \frac{A_1}{r^2} + \frac{2A_1^2}{c^2r^3} - \frac{4A_1^2}{3c^2r^4} + \dots \quad (29)$$

Let us seek the solution of the interior field equation (22) as

$$f^-(r) = f_c^-(r) + f_p^-(r) \quad (30)$$

where, f_c and f_p are complementary and particular solutions, respectively.

The complementary solution of the interior field equation

$$f_c^-(r) = B_0 \quad (31)$$

where, B_0 is an arbitrary constant.

Seeking the solution in the form

$$f_p^-(r) = D_2r^2 + D_3r^3 + D_4r^4 + \dots \quad (32)$$

Differentiating equation (32) with respect to r yields

$$f_p'^- = 2D_2r + 3D_3r^2 + 4D_4r^3 + \dots \quad (33)$$

$$f_p''^- = 2D_2 + 6D_3r + 12D_4r^2 + \dots \quad (34)$$

where D_2 , D_3 and D_4 are arbitrary constants

Putting equation (32), (33) and (34) into (21) and equating coefficients of r^0 , r^1 and r^2 on both sides and using the well - known physical relationship between mass and density for a sphere of radius R we obtain equations which gives

$$D_2 = \frac{GM}{2R^3} \quad (35)$$

$$D_4 = -\frac{3G^2M^2}{8c^2R^6} \quad (36)$$

The general solution of the interior field equation (30) is given by

$$f^-(r) = B_0 + D_2r^2 + D_3r^3 + D_4r^4 + \dots \quad (37)$$

By the condition of continuity of gravitational scalar potential function across boundaries, $r = R$

$$B_0 + D_2R^2 + D_3R^3 + D_4R^4 + \dots = \frac{A_1}{r} + \frac{A_2}{r^2} + \frac{A_3}{r^3} + \dots \quad (38)$$

By the condition of continuity of normal derivatives of gravitational scalar potential function across all boundaries, $\left(\frac{\partial F^+}{\partial r}\right)_{r=R} = \left(\frac{\partial F^-}{\partial r}\right)_{r=R}$

$$2D_2R + 4D_4R^3 + \dots = -\frac{A_1}{R^2} - \frac{2A_2}{R^3} - \frac{3A_3}{R^4} + \dots \quad (39)$$

Solving equation (39) for A_1 and using the well - known physical relationship between mass and density for a sphere of radius R gives

$$A_1 = GM \left(1 - \frac{GM}{c^2R} + \frac{GM}{c^2R^2} + \dots\right) \quad (40)$$

$$A_4 = -\frac{4}{3c^2}A_1^2 = -\frac{4G^2M^2}{3c^2} + \dots \quad (41)$$

Solving equation (38) for B_0 to the order of c^{-2} gives

$$B_0 = \frac{GM}{2R} + \frac{GM}{R^2} + \frac{9G^2M^2}{8c^2R^2} + \frac{G^2M^2}{2c^2R^3} + \dots \quad (42)$$

Substituting equation (40) and (41) into the exterior field equation (22) gives

$$f^+(r) = \frac{GM}{r} \left\{1 - \frac{3GM}{2c^2R}\right\} + \frac{GM}{r^2} \left\{1 + \frac{3GM}{2c^2R}\right\} + \frac{GM}{r^3} \left\{1 + \frac{3GM}{2c^2R}\right\} - \frac{4G^2M^2}{3c^2r^4} + \dots \quad (43)$$

Similarly, substituting equations (35), (36) and (42) into (37) gives the interior field equation as

$$f^-(r) = B_0 + \frac{3GM}{2R}r^2 - \frac{3G^2M^2}{8c^2R^6}r^4 + \dots \quad (44)$$

where $k = GM$, M is the rest mass of the sun, G is the universal gravitational constant, c is the speed of light in vacuum, R is the radius of the spherical astrophysical body and $r > R$ for the exterior field, $f^+(r)$ and $f^-(r)$

are the extended Newtonian gravitational scalar potential exterior and interior.

Equations (43) and (44) are the extended exterior and interior Newtonian gravitational scalar potential based on the golden Riemannian geometry of space - time. It is interesting to note that these equations reduce in the limit of c^0 to the corresponding pure Newtonian expression and to the order of c^{-2} it contains post - Newton and post - Einstein correction terms. It is interesting to note that these results contain additional terms not found in [11]. The immediate consequences of the additional correction terms in the extended Newtonian gravitational scalar exterior and interior are that they can be:

- Substituted into Newton's, Einstein's and General equations of motion to obtain generalizations of Newton's, Einstein's and General equations of motion to obtain corresponding revisions to the planetary equation of motion and hence the planetary parameter such the anomalous orbital precession in the solar, gravitational shift, orbital eccentricity, orbital period, orbital amplitude, perihelion and aphelion distance.
- Applied to examine the theoretical analysis of the motion of particles of non-zero rest masses, gravitational length contraction and time dilation and in simple linear harmonic oscillator to determine its energy levels.
- Applied to the motion of test particles in the gravitational fields of static homogeneous spherical bodies to give more accurate expression for the motion of test particles as it introduces hitherto unknown correction terms.
- Can be substituted into the well - known Einstein's (geodesic) equation of motion and the resulting equations of motion solved to obtain the corresponding post -Einstein and Newtonian generalizations and the physical effects.

4. Conclusions

We have in this work shown one possible way to construct an extended Newtonian dynamical gravitational field equation and scalar potential exterior and interior for static homogeneous spherical massive bodies using the golden Riemannian geometry of space - time. It is interesting and instructive to note that these equations reduce in the limit of c^0 , to the corresponding pure Newtonian terms implying that it agrees with the well - known equivalence principles in Physics and to the order of c^{-2} it contains unknown post - Newton and post - Einstein correction terms. The door is

therefore open for the development and experimental investigations and possible applications of this theory.

ACKNOWLEDGEMENTS

The authors wish to sincerely appreciate the effort of **Tertiary Education Trust Fund (TETFUND), 2019/2020**, for funding this research work from inception to the final stage.

Source (s) of Funding

This research was funded by the Tertiary Education Trust Fund under the Institution Based Research grant.

REFERENCES

- [1] Einstein A. Relativity: the special and the general theory. General Press, 2016.
- [2] Anderson A. J. Principles of relativity physics, acad. Press, New York and London, 1967.
- [3] S. X. K. Howusu. Complete dynamical theory of physics. Jos University Press Ltd, pages 31–35, 2010.
- [4] S. X. K. Howusu. Riemannian revolution in physics and mathematics. Jos University Press Ltd, pages 1 – 190, 2013.
- [5] Foster J. A. and Nightingale J. D. A short course in General Relativity. Springer Science & Business Media, 2010.
- [6] Schwarzschild K. On the gravitational field of a point-mass, according to Einstein's theory. American Association of Physics Teachers, page 19, 1982.
- [7] L.W. Lumbi, S. X. K. Howusu and M. S. Liman. Generalized dynamical gravitational field equation for static homogeneous spherical massive distribution of mass. International Journal of Modern Theoretical Physics 3 (1): 37, 43, 2014.
- [8] Y. Nura, S. X. K. Howusu, A. Musa N. Ibrahim, and A. Babakura. The general equation of motion in a gravitational field based upon the golden metric tensor. American Journal of Modern Physics, 6(6): 127–131, 2017.
- [9] Bergman P. G. Introduction to the theory of Relativity. Prentice Hall, India, 1987.
- [10] Weinberg S. Gravitation and cosmology: Principles and applications of general theory of relativity, copyright 1972 by John Wiley & Sons, 1972.
- [11] M.R. Spiegel and L. J. Stephens. Schaum's Outline of Theory and Problems of. McGraw, 1961.