

The Higgs-like Bosons and Quark Compositeness

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Abstract Considering that each quark is composed of two prequarks, called primons, it is shown that the recently found neutral Higgs-like boson belongs to a triplet constituted of a neutral boson H^0 and two charged bosons H^+ and H^- , and that H^0 is, actually, a triplet and both H^+ and H^- are doublets. The quantum numbers of these bosons are calculated and shown to be associated to a new kind of hypercharge which is directly related to the weak decays of hadrons and to the CKM matrix elements. Solutions to the proton spin puzzle and to other problems of particle physics are presented.

Keywords Higgs Boson, Higgs-like Bosons, CKM Matrix, Weak Decays, Preons, Prequarks, Primons

1. Introduction

The collaborations ATLAS[1,2] and CMS[3] have recently discovered a light narrow resonance around 126 GeV. This remarkable discovery, supported by evidence from the Tevatron[4], has triggered discussions on the nature of such resonance because both ATLAS and CMS report discrepancies with respect to the Standard Model (SM) backgrounds. The data have established that it is indeed a boson, probably neutral with spin zero. The identification of the new resonance with the long-sought Higgs boson of the Standard Model is not certain at all. In particular, D0 and CDF have reported an excess of the top quark forward-backward asymmetry (A_{FB}^t) [5,6] and an enhancement in the $Vb\bar{b}$ channel[4], while the LHC has reported an excess in the diphoton channel[7,8]. These disagreements in the Higgs couplings, as pointed out by Choudhury *et al.*[9], is “a Window to New Physics”.

Because of this possibility there has lately been an array of different proposals offering alternatives for new physics and treating the newly found resonance as a Higgs-like boson. There are many recent references in this line such as, for example, the eight references[10,11,12,13,14,15,16,17]. One sees that there are works proposing Higgs doublets and Higgs triplets. In particular, Low, Lykken & Shaughnessy[17] propose that the newly found boson is the neutral member of an electroweak triplet of a custodial 5-plet.

This work presents a different view from all the works above mentioned and brings to the discussion the possibility of a composite quark. Although prequarks (or preons) have not yet been directly proven by any experiment, there are

several indirect evidences of their existence. Some examples are: **1)** the distributions of charge in nucleons, found by Hofstadter & Herman[18] that cannot be explained by three point-like quarks; **2)** the *ad hoc* Kobayashi-Maskawa matrix elements; **3)** the so-called EMC effect[19,20] according to which quarks of nucleons are slower when nucleons are inside nuclei; **4)** the SLAC E143 Collaboration[21], the Spin Muon Collaboration[22] and Hermes Collaboration[23] have found that the three quarks of the proton account for only half of its total spin which is a fact called the proton spin puzzle; **5)** G. Miller at Argonne[24] found that close to its center the neutron has a negative charge equal to $-1/3e$ (inside the positive region with $+1/2e$). Another important effect was discovered by the group of Prof. Alan Krusch that conducted the interaction of a polarized proton beam on a polarized proton target. The data show[25] that, when the collision intensity is higher, protons with parallel spins interact much stronger than protons with anti-parallel spins. According to QCD both cases would have the same strength. Let us recall that according to QCD's asymptotic freedom the forces between quarks become weaker at short distances.

My argumentation is that QCD is not wrong. It cannot explain the Krusch effect because it is due to the substructure of quarks. I present in section 7 an explanation for this effect and for the other ones above mentioned.

Commenting on the discovery of quarks at SLAC, M. Riordan[26] says “A way to interpret this unexpected behavior was that the electrons were hitting some kind of hard core inside the target protons”. Taking into account the quark model of today we know that such core was not made of quarks. Later on, below, we will grasp the true nature of this core.

Ending the introduction I can say that preons have not yet been directly found by the LHC because they should be very light fermions and, thus, they are very hard to be directly seen, and thus with the LHC energies we only see the bosons

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related to their interaction. As it will be shown below there should exist a neutral boson (already found by the LHC) and two charged bosons (yet to be found).

2. The Quark Compositeness Model

Quark compositeness is an old idea, proposed since the 1970s[27-37]. In order to distinguish the model of this work from the other models found in the literature above cited I name the prequarks of this model with a different name: primons. Primon is a word derived from the Latin word primus which means first.

A major distinction is that in this work leptons are supposed to be elementary particles. This is actually consistent with the smallness of the electron mass which is already too small for a particle with a very small radius[38]. Another important distinction is that this work breaks the quark into two particles instead of three particles. Some major parts of the model have been presented in references[39,40].

2.1. Compositeness of Quarks Flavors and Colors

In order to reproduce the spectrum of 6 quarks and their colors we need 4 primons (p_1, p_2, p_3, p_4) in 3 supercolor states (α, β, γ). Each color is formed by the two supercolors of the two primons forming a particular quark (Table 1). As to charge, primons have charges $+(5/6)e$ and $(-1/6)e$ according to Table 2.

With this construction of quarks flavors we solve one of the problems of particle physics: there are only six quarks. Of course, there is a similar construction for the antiquarks (with antiprimons). Using Table 2 we obtain Table 3 for the flavors of quarks.

Table 1. Generation of colors from supercolors

	α	β	γ
α		blue	green
β	blue		red
γ	green	red	

Table 2. Electric charges of primons

Superflavor	Charge
p_1	$+5/6$
p_2	$-1/6$
p_3	$-1/6$

p_4	$-1/6$
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Table 3. Composition of quark flavors

	p_1	p_2	p_3	p_4
p_1		u	c	t
p_2	u		d	s
p_3	c	d		b
p_4	t	s	b	

2.2. The New Hypercharge and the New SU(2)

For primons we should have the baryon number $B=1/6$, of course. As it is described in references[39,40], going on the footsteps of Gell-Mann and Nishijima (for the quark model), we can use the formulas

$$Q = 2B + \frac{1}{2}(P_1 + P_2 + P_3 + P_4) \quad (1)$$

for a system of primons (a quark), where, $P_1 = +1$, for primon p_1 , and $P_j = -1$ ($j = 2, 3, 4$) for the other primons. Thus, we have

$$\frac{2}{3} = 2 \times \left(\frac{1}{6} + \frac{1}{6} \right) + \frac{1}{2}(1 + (-1)) \quad (2)$$

for quarks u, c, t , and

$$-\frac{1}{3} = 2 \times \left(\frac{1}{6} + \frac{1}{6} \right) + \frac{1}{2}(-1 + (-1)) \quad (3)$$

for quarks d, s, b .

As quarks u and d have isospins equal to $+1/2$ and $-1/2$, respectively, we are forced to have $I_3 = \pm 1/4$ for primons, p_1, p_2 and p_3 . Making use of the modified Gell-Mann—Nishijima relation

$$Q = I_3 + \frac{1}{2}(B + \Sigma_3) \quad (4)$$

for primons, we obtain

$$\frac{5}{6} = \frac{1}{4} + \frac{1}{2} \left(\frac{1}{6} + 1 \right) \quad (5)$$

for primon p_1 , and

$$\begin{aligned} -\frac{1}{6} &= \frac{1}{4} + \frac{1}{2} \left(\frac{1}{6} + (-1) \right) \\ -\frac{1}{6} &= -\frac{1}{4} + \frac{1}{2} \left(\frac{1}{6} + 0 \right) \end{aligned} \quad (6)$$

for p_2, p_3 and p_4 . Therefore, Σ_3 can assume the values $-1, 0$ and $+1$ and, thus, these values can be considered as the projections of $\Sigma = 1$ (Table 4).

In terms of Σ and Σ_3 a quark has the four possible states $|\Sigma, \Sigma_3\rangle$ (described in detail in ref.[39,40]):

$$\begin{array}{|c|} \hline c, t : |1, +1\rangle \\ d : |1, 0\rangle \quad u : |0, 0\rangle \\ s, b : |1, -1\rangle \\ \hline \end{array}$$

Table 4. Projections of the isospin I and of the quantum number Σ

	I_3	Σ_3
p_1	$+\frac{1}{4}$	$+1$
p_j	$+\frac{1}{4}$	-1
$(j = 2, 3, 4)$	$-\frac{1}{4}$	0

With the above assignments, the quark doublets

$$\begin{pmatrix} u \\ d \end{pmatrix}, \begin{pmatrix} c \\ s \end{pmatrix}, \begin{pmatrix} t \\ b \end{pmatrix}$$

are written as

$$\begin{pmatrix} |0, 0\rangle \\ |1, 0\rangle \end{pmatrix}, \begin{pmatrix} |1, +1\rangle \\ |1, -1\rangle \end{pmatrix}, \begin{pmatrix} |1, +1\rangle \\ |1, -1\rangle \end{pmatrix}.$$

Organizing the values of I_3 and Σ_3 for quarks on a table we obtain Table 5 below.

Table 5. Assignments of the values of I_3 and Σ_3 for quarks

	I_3	Σ_3
c, t	0	$+1$
u	$+1/2$	0
d	$-1/2$	0
s, b	0	-1

We can also construct the diagram below (Fig 1) that shows how the decays of quarks are related to the new hypercharge and to isospin. Of course, there is a close relation between the diagram and the CKM matrix elements.

Comparing Fig. 1 with the Kobayashi-Maskawa matrix[41] (with 3 decimals)

$$\begin{pmatrix} |U_{ud}| & |U_{us}| & |U_{ub}| \\ |U_{cd}| & |U_{cs}| & |U_{cb}| \\ |U_{td}| & |U_{ts}| & |U_{tb}| \end{pmatrix} = \begin{pmatrix} 0.974 & 0.225 & 0.004 \\ 0.225 & 0.974 & 0.041 \\ 0.009 & 0.040 & 0.999 \end{pmatrix}$$

we notice that the matrix elements $|U_{cs}|$ and $|U_{tb}|$, which are about 1, satisfy the selection rule $\Delta\Sigma_3 = -2$, $\Delta\Sigma = 0$, $\Delta I_3 = 0$. The other very large element $|U_{ud}|$ (which is also close to 1) satisfies the selection rule $\Delta\Sigma_3 = 0$, $\Delta\Sigma = +1$, $\Delta I_3 = +1$. From these three matrix elements we obtain the rule: The largest elements have either the maximum $|\Delta\Sigma_3|$ or the maximum $|\Delta I_3|$ within the same quark doublet. The other large elements $|U_{cd}| = 0.225$ and $|U_{us}| = 0.225$ obey, respectively, the selection rule $|\Delta\Sigma_3| = 1$, $\Delta\Sigma = 0$, and $|\Delta I_3| = 1/2$. From the diagram we see why $|U_{cd}| \approx |U_{us}|$. According to the diagram we should have $|U_{ts}| \approx |U_{cb}|$ and $|U_{td}| \approx |U_{ub}|$ and this is indeed the case.

Therefore, the CKM matrix elements are directly related to the substructure of quarks. Let us recall that within the Standard Model the CKM matrix is an *ad hoc* element.

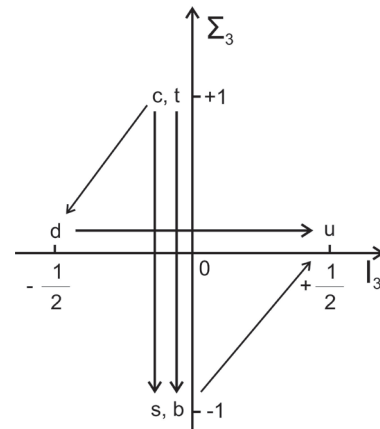


Figure 1. Diagram showing how weak decays are related to the variations of Σ_3 and I_3

3. The Weak Decays of Hadrons

Taking into account the values of the new quantum numbers Σ_3 for quarks, de Souza[42] has shown that the weak decays of all hadrons (leptonic, semileptonic and nonleptonic) are directly related to the variations of Σ_3 , that is, to $\Delta\Sigma_3$, and, thus, there is a vectorial current j_μ^Σ between quarks. As an example, I present on Table 6 the semileptonic decays of light baryons.

4. The Masses of Primons

The magnetic moments of primons should be given by $\mu_1 = \frac{5}{6} \frac{e}{2m_1}$ for p_1 and $\mu_2 = -\frac{1}{6} \frac{e}{2m_2}$ for p_2, p_3, p_4 and, hence $\mu_1 = -\frac{5m_2}{m_1} \mu_2$, $\mu_3 = \frac{m_2}{m_3} \mu_2$. Since the spin content of quarks should be the same we have

$$\frac{\mu_u}{\mu_d} = \frac{\mu_1 + \mu_2}{\mu_2 + \mu_3} \quad (7)$$

and since $\mu_u / \mu_d = -2$, and using the above relations we obtain

$$-2 = -\frac{5m_2}{m_1} \left(\frac{1 - \frac{m_1}{5m_2}}{1 + \frac{m_2}{m_3}} \right). \quad (8)$$

Making $m_3 = fm_2$ and solving for the ratio m_1/m_2 we arrive at

$$\frac{m_1}{m_2} = \frac{5}{3 + \frac{2}{f}} \quad (9)$$

Table 6. Semileptonic decays of light baryons with their corresponding Cabibbo factors and $\Delta\Sigma_3$

Decay	Cabibbo factor	$\Delta\Sigma_3$
$n \rightarrow p$	$\cos \theta_c$	0
$\Sigma^+ \rightarrow \Lambda$	$\cos \theta_c$	0
$\Sigma^- \rightarrow \Lambda$	$\cos \theta_c$	0
$\Sigma^- \rightarrow \Sigma^0$	$\cos \theta_c$	0
$\Lambda \rightarrow p$	$\sin \theta_c$	+1
$\Sigma^- \rightarrow n$	$\sin \theta_c$	+1
$\Xi^- \rightarrow \Lambda$	$\sin \theta_c$	+1
$\Xi^- \rightarrow \Sigma^0$	$\sin \theta_c$	+1
$\Xi^0 \rightarrow \Sigma^+$	$\sin \theta_c$	+1
$\Xi^- \rightarrow \Xi^0$	$\cos \theta_c$	0

and as the mass of $u(p_1 p_2)$ and $d(p_2 p_3)$ are approximately equal, it is reasonable to consider $m_3 \approx m_1$. Supposing that $m_2 \approx m_1$, we obtain $f \approx 1$ which is very consistent. Therefore, primons should have approximately the same mass which is an inherent mass such as that of leptons.

5. The Arrangement of Primons in the Nucleons and the Stability of the Proton

5.1. The Arrangement of Primons in the Nucleons

Taking into account the results of Hofstadter & Hermann[18] on the charges of the nucleons we can say that both nucleons should have an inner common central core constituted of primons p_1, p_2, p_3 . The outer primons of the proton are p_1, p_1, p_2 according to Fig. 2 below. The outer primons of the neutron are p_2, p_2, p_3 , forming quarks u, d, d between the two layers.

It is worth mentioning that this picture of the nucleons completely agrees with the results of Miller[24] who found a charge of $-1/3$ inside the central positive core of the neutron. This $-1/3$ charge is the result of $-1/3 = -1/6 + (-1/6)$ of primons p_2 and p_3 . This picture of the nucleons also completely agrees with the results of Povn and Hüfner[43] that have found an average size of about 0.5 fm for constituent quarks u and d .

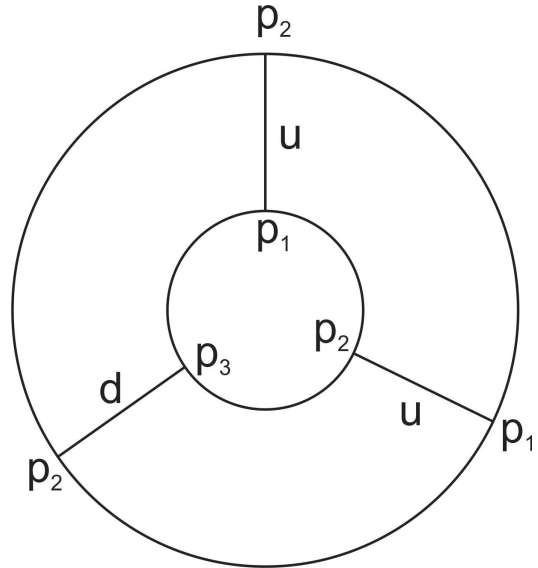


Figure 2. Pictorial representation of the arrangement of primons in the proton

5.2. The Stability of the Free Proton

Taking into account the stability of the (free) proton we can deduce that p_1 and p_2 should be stable against weak decay. What about p_3 ? From the instability of the free neutron we deduce that p_3 should suffer weak decay, but

only the outermost p_3 because the innermost p_3 is also present in the proton. We can understand why this happens taking a look at the decays of hadrons and observing that weak decays do not occur when both $\Delta\Sigma_3 = 0$ and $\Delta I_3 = 0$.

Let us recall that $u = p_1 p_2$ so that $\Sigma_3(u) = \Sigma_3(p_1) + \Sigma_3(p_2) = +1 - 1 = 0$, and that $\Sigma_3(d) = \Sigma_3(p_2) + \Sigma_3(p_3) = 0 + 0 = 0$. As to I_3 we have that $I_3(u) = I_3(p_1) + I_3(p_2) = +1/4 + 1/4 = +1/2$, and $I_3(d) = I_3(p_2) + I_3(p_3) = -1/4 - 1/4 = -1/2$. Therefore, the inner layer of the proton (p_1, p_2, p_3) has $\Sigma_3 = 1 + (-1) + 0 = 0$, $I_3 = +1/4 + 1/4 - 1/4 = +1/4$ because the p_2 of the inner layer belongs to u . In the outer layer (p_2, p_1, p_2) the first p_2 belongs to u and the other one belongs to d , so that $\Sigma_3 = -1 + 1 + 0 = 0$. As to I_3 we have $I_3 = +1/4 + 1/4 + (-1/4) = +1/4$, and thus $\Delta\Sigma_3 = 0$, $\Delta I_3 = 0$ between the two layers. Therefore, the free proton is always stable. For the neutron the inner layer is the same as in the proton, and the outer layer (p_2, p_2, p_3) has $\Sigma_3 = -1 + 0 + 0 = -1$, $I_3 = 3 \times (-1/4) = -3/4$, and thus between the innermost and outermost layers $|\Delta\Sigma_3| = 0 - (-1) = 1$, $\Delta I_3 = 1/4 - (-3/4) = 1$, and thus, the free neutron always decays.

6. The Higgs-like Bosons and their Quantum Numbers

6.1. The Spins of Primons and of the Higgs-like Bosons

In order to have $S_z = \pm 1/2$ for quarks, we should have $S_z = \pm 1/4$ for primons. Thus, we postulate that primons are fermions ($S = 1/2$), but with $S_z = \pm 1/4$. Taking a look at Fig. 2 each layer of primons should have $\pm \frac{1}{4} + \left(\pm \frac{1}{4}\right) + \left(\mp \frac{1}{4}\right) = \pm \frac{1}{4}$, so that the two layers together should always have $S_z = \pm 1/2$, of course. Therefore, the Higgs-like bosons should have $S = 0$, and are, thus, scalar bosons. On subsection 7.3 on the solution of the Proton Spin Puzzle we return to discuss the spins of primons.

6.2. The Quantum Numbers Σ_3 of the Higgs-like Bosons

Taking into account the charges of primons from Table 2, we see that there should be three Higgs-like bosons, H^0, H^+, H^- for generating the masses of quarks as it is shown on Table 7, because

$$+\frac{5}{6} - \left(-\frac{1}{6}\right) = +1; -\frac{1}{6} - \left(+\frac{5}{6}\right) = -1, \text{ and } -\frac{1}{6} - \left(-\frac{1}{6}\right) = 0.$$

Let us now assign the quantum numbers Σ_3 for the Higgs-like bosons. As $u = p_1 p_2$, its Σ_3 is $\Sigma_3(u) = \Sigma_3(p_1) + \Sigma_3(p_2) = +1 + (-1) = 0$, and so we have $|\Sigma_3(H^+)| = |\Sigma_3(H^-)| = |\Sigma_3(p_2) - \Sigma_3(p_1)| = 2$. For d we have $\Sigma_3(d) = \Sigma_3(p_2) + \Sigma_3(p_3) = 0 + 0 = 0$, and thus, $|\Sigma_3(H^0)| = 0$. Doing the same for the other quarks, and summarizing the results on a table, we obtain Table 8 below.

Table 7. The generators of quark masses

Quark	Mass (GeV)	Charge	Higgs-like boson
$u(p_1 p_2)$	0.3	$+2/3$	H^+, H^-
$c(p_1 p_3)$	1.5	$+2/3$	H^+, H^-
$t(p_1 p_4)$	170	$+2/3$	H^+, H^-
$d(p_2 p_3)$	0.3	$-1/3$	H^0
$s(p_2 p_4)$	0.5	$-1/3$	H^0
$b(p_3 p_4)$	4.5	$-1/3$	H^0

The three Higgs-like bosons above mentioned were predicted by de Souza in reference[40] and are implicit in reference[39].

6.3. The Masses of the Higgs-like Bosons

The mass of H^0 has been measured by the collaborations ATLAS[1,2] and CMS[3] as being about 126 GeV. This corresponds to a Compton wavelength around 10^{-18} m. It is important to point out that the mass of H^0 is of the same order of that of the weak boson Z^0 . The charged weak bosons W^+, W^- have a mass $M(W^\pm) = M(Z^0) - 10 \text{ GeV}$. I think it is not purely coincidental that the mass of H^0 and of the weak bosons have the same order of magnitude. From this we may guess that the mass of H^+ and H^- is $M(H^\pm) = M(H^0) - 10 \text{ GeV} = 116 \text{ GeV}$.

Table 8. The quantum numbers of the Higgs-like bosons

	H^0	H^+, H^-
Σ_3	$\pm 1(s, b)$	$\pm 2(u)$
	$0(d)$	$\pm 1(c, t)$

As it is shown on Table 8, there are 3 H^0 which probably have 2 slightly different masses, that is, the H^0 s with $\Sigma_3 = \pm 1$ may have a different mass from the H^0 with $\Sigma_3 = 0$. In fact, CMS and ATLAS have reported for the Higgs boson the respective masses of (125.8 ± 0.7) GeV and (126.8 ± 0.9) GeV[44] which may be the two different masses above mentioned. The same may hold for the charged bosons. The H^+ with $\Sigma_3 = \pm 2$ and the H^+ with $\Sigma_3 = \pm 1$ may have slightly different masses. The H^- bosons, of course, should have the same behavior.

7. The Higgs-like Lagrangian

In order to construct a gauge theory for this new interaction we should consider a Lagrangian for three massive spin-0 bosons, of the type

$$L = \frac{1}{2} \partial_\mu \phi_+ \partial^\mu \phi_+ + \frac{1}{2} \partial_\mu \phi_- \partial^\mu \phi_- + \frac{1}{2} \partial_\mu \phi_0 \partial^\mu \phi_0 - \frac{1}{2} m_+ \phi_+^2 - \frac{1}{2} m_- \phi_-^2 - \frac{1}{2} m_0 \phi_0^2 \quad (10)$$

where ϕ_+, ϕ_-, ϕ_0 refer to the three Higgs-like fields corresponding to H^+, H^-, H^0 , respectively, with their corresponding masses m_+, m_-, m_0 . This is a well known Lagrangian found in the literature which satisfies Klein-Gordon equation. For example, in reference[45] there is a detailed discussion on this subject. We just have to consider that, according to what was above discussed, we have a $(2 \times 1 + 1)$ -dimensional $SU(2)$ directly linked to Σ . That is, there are the $SU(2)$ bosonic current

$$j_b^{\mu\Sigma} = i \left[\left(\phi^\Sigma \right)^\dagger T^\Sigma \partial^\mu \phi^\Sigma - \left(\partial^\mu \phi^\Sigma \right)^\dagger T^\Sigma \phi^\Sigma \right] \quad (11)$$

and the $SU(2)$ fermionic current

$$j_f^{\mu\Sigma} = \left(\bar{\psi}^\Sigma \right) \gamma^\mu T^\Sigma \psi^\Sigma. \quad (12)$$

In these currents T^Σ are matrices of the $(2 \times 1 + 1)$ -dimensional representation of $SU(2)$.

Therefore, we should modify the Standard Model replacing the Higgs field Lagrangian by the above Higgs-like field Lagrangian.

For the interaction between the two primons that form a particular quark the Lagrangian is much more complex. One of the complications comes from the fact of having $S_z = \pm 1/4$ for each primon. Since Dirac equation has the

Pauli spin matrices included inside the γ^μ matrices, we should have a modified Dirac equation for primons. Another complication comes from the fact that the spins of the primons of a quark should change together when, for example, a quark interacts with a gluon since gluons have

spin 1. This means that for a pair of primons forming a quark

$$[\vec{S}_1, \vec{S}_2] \neq 0$$

and thus the spin degrees of freedom are not independent from each other. This leads to some interesting properties regarding the nucleon as a whole as discussed in reference[39].

Thus, the Lagrangians of the two primons that form a quark have to be coupled. We should investigate this matter much further.

8. Solutions to Some Important Problems of Particle Physics

8.1. The Electric Charge Densities in the Nucleons

The results of the experiment of Hofstadter & Hermann are the starting point for the model. We notice that there is a complete coherence between the experimental results and the model. And moreover, the model completely agrees with the findings of Miller[24] according to whom there is a negative charge of $-1/3$ inside the positive central part of the charge density of the neutron. Neither result can be accounted for by three point-like quarks.

8.2. The Number of Quarks

The 4 primons arranged in pairs with different flavours form exactly 6 quarks.

8.3. The Proton Spin Puzzle

As it was shown by de Souza in references[40] and[42], the puzzle is easily solved. Here we just mention the main points. Deep Inelastic Scattering (DIS) experiments have probed the inner layer of primons of the proton, and that is why the results yield about half of the total spin. Taking into account the other layer of primons we arrive at a total spin $S_z = +1/2$ (or $S_z = -1/2$) as it should be.

8.4. The CKM Matrix Elements and Cabibbo Factors

The results of section 2.2 above and the results of reference[42] clearly pave the way for the understanding of the CKM matrix elements which are presently completely *ad hoc*. They are directly related to the variations of Σ_3 in weak decays. As it is seen in reference[42] Cabibbo factors are also directly related to the variations of Σ_3 in weak decays.

8.5. The Dual Roles of Quarks: Valence and Constituent Quarks

Valence quarks are, actually, primons, and are almost massless, and constituent quarks are the true quarks. Thus, there is a confusion in the interpretation of the data. The three outermost primons have not yet been directly seen because they become blurred into the quark sea in deep inelastic scattering experiments and only the three innermost primons have been seen together as an inner core and two of them

have been seen together according to G. Miller findings (my interpretation). This is the same inner core that was seen by Hofstadter & Hermann and was also reported by M. Riordan[26] on the description of the discovery of quarks. This subject is discussed in more detail in references[40] and[42].

8.6. The EMC Effect

The EMC effect is due to the interactions between primons of the outermost layer of a nucleon and primons of the outermost layer of another close nucleon. That is, there should be something like primon bonds between two close nucleons. The model opens up the possibility for the understanding of the hole inside the alpha particle and its point group symmetry which is C_{3v} [46], that is, it has a pyramidal shape. With point-like quarks we will never understand how such arrangement of nucleons such as the deuteron and the alpha particle can be formed.

8.7. The Krisch Effect

As it was shown in the introduction, Prof. Alan Krisch group conducted the interaction of a polarized proton beam on a polarized proton target and found[25] that, at higher collision intensities, protons with parallel spins interact stronger than protons with anti-parallel spins.

At very high energies, most interactions arise from the interactions between the two central layers of primons of the two protons, that is, between the set (p_1, p_2, p_3) of a proton and the set (p_1, p_2, p_3) of the other proton.

When the two protons have parallel spins, in terms of spins we have for each central layer of primons two spins up and one spin down (or the other way around) and, thus, there are, 04 interaction with $(\uparrow\uparrow)$ and 01 with $(\downarrow\downarrow)$ (for two spins up, and one spin down) because the Higgs-like bosons have $S=0$. When the proton spins are antiparallel we have two spins up and one spin down for one proton, and two spins down and one spin up for the other proton (or the other way around). In this case we have 02 interactions with $(\uparrow\uparrow)$ and 02 interaction with $(\downarrow\downarrow)$. Therefore, there is a ratio of $5/4 = 1.25$ between the two cases. That is, there are about 11.12% more transitions for parallel protons than for antiparallel protons.

8.8. The Masses of Baryons and the Harmonic Approximation for Effective Potentials

Taking a look at Fig. 2 we observe that when protons collide at high energies and form baryon resonances, the quarks oscillate quite a lot and can be considered as quantum oscillators. That is why de Souza[47] was able to describe the masses of almost all baryons using such an approximation. The first excited state of Λ_b^0 , which had been predicted to have a mass of 5.93 GeV, has recently been reported[48] by CDF with a mass of about 5.92 GeV, just

0.17% off.

9. Conclusions

I have presented a completely new view for the Higgs-like bosons in terms of a substructure for quarks. It is a very coherent model, based on experimental data. The quantum numbers of the Higgs-like bosons are presented and are directly related to the weak decays of hadrons. Solutions have been presented to some important problems of the Physics of hadrons and quarks. It is a model that can be tested further with more data from the LHC.

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REFERENCES

- [1] G. Aad *et al.* (ATLAS Collaboration), Phys. Lett. B, vol. 716, 1, 2012.
- [2] G. Aad *et al.* (ATLAS Collaboration), Phys. Rev. D, vol 86, 032003, 2012.
- [3] S. Chatrchyan *et al.* (CMS Collaboration), Phys. Lett. B, vol. 716, 30, 2012.
- [4] T. Aaltonen *et al.* (CDF and D0 Collaborations), Phys. Rev. Lett., vol. 109, 071804, 2012.
- [5] T. Aaltonen *et al.* (CDF and D0 Collaborations), CDF note 10807, 2011.
- [6] V. M. Abazov *et al.* (D0 Collaboration), Phys. Rev. D, Vol. 84, 112005, 2011.
- [7] ATLAS Collaboration, ATLAS-CONF-2012-168.
- [8] CMS Collaboration, CMS-PAS-HIG-12-015.
- [9] D. Choudhury, R. Islam, A. Kundu, and B. Mukhopadhyaya, "Anomalous Higgs Couplings as a Window to New Physics", arXiv 1212.4652v1, 19 Dec., 2012.
- [10] M. Carena *et al.*, JHEP, vol. 1203, 14, 2012.
- [11] M. Carena *et al.*, JHEP, vol. 1207, 175, 2012.
- [12] J. Cao *et al.*, JHEP, 1210, 79, 2012.
- [13] M. Carena, I. Low and C. E. M. Wagner, JHEP, vol. 1208, 60, 2012.
- [14] A. G. Akeroyd and S. Moretti, Phys. Rev. D, vol 86, 035015, 2012.
- [15] M. B. Voloshin, Phys. Rev. D, vol 86, 093016, 2012.
- [16] A. Alves *et al.*, Phys. Rev. D, vol 84, 115004, 2011.

- [17] I. Low, J. Lykken and C. Shaughnessy, arXiv 1207.1093v1, 4 Jul., 2012.
- [18] R. Hofstadter and R. Herman, Phys. Rev. Lett., vol. 6, 293, 1961.
- [19] J. Ashman *et al.* (European Muon collaboration), “An Investigation of the Spin Structure of the Proton in Deep Inelastic Scattering of Polarized Muons on Polarized Protons”, Nucl. Phys. B, vol. 328, 1, 1989.
- [20] J. Seely *et al.*, New Measurements of the European Muon Collaboration Effect in Very Light Nuclei”, Phys. Rev. Lett., vol. 103, 202301, 2009.
- [21] K. Abe *et al.* (SLAC E143 Collaboration), “Measurements of the Proton and Deuteron Spin Structure Functions g_1 and g_2 ”, Phys. Rev. D, vol. 58, 112003, 1998.
- [22] B. Adeva *et al.* (Spin Muon Collaboration), “Spin Asymmetries $A(1)$ and Structure functions g_1 of the Proton and the Deuteron from Polarized High energy Muon Scattering”, Phys. Rev. D, vol. 58, 11201, 1998.
- [23] A. Airapetian *et al.* (Hermes Collaboration), “Precise determination of the spin structure function g_1 of the proton”, Phys. Rev. D, vol. 75, 012007, 2007.
- [24] G. A. Miller, “Charge Densities of the Neutron and Proton”, Phys. Rev. Lett., vol. 99, 112001, 2007.
- [25] A. D. Krisch, “Hard collisions of spinning protons: Past, present and future”, Eur. Phys. J. A, vol. 31, 417, 2007.
- [26] M. Riordan, “The Discovery of Quarks”, SLAC-PUB-5724, April 1992.
- [27] J. C. Pati, A. Salam, “Lepton number as the fourth “color””, Phys. Rev. D, vol. 10, 275-289, 1974, erratum published as Phys. Rev. D, vol. 11(3), 703, 1975.
- [28] H. Terazawa, Y. Chikashige and K. Akama, “Unified model of the Nambu-Jona-Lasinio type for all elementary particles”, Phys. Rev. D, vol. 15(2), 480, 1977.
- [29] Y. Ne’eman, “Irreducible gauge theory of a consolidated Weinberg-Salam model”, Phys. Lett. B, vol. 81(2), 190, 1979.
- [30] H. Harari, “A schematic model of quarks and leptons”, Phys. Lett. B, vol. 86, 83, 1979.
- [31] M. A. Shupe, “A composite model of leptons and quarks”, Phys. Lett. B, vol. 86, 87, 1979.
- [32] H. Fritzsch, G. Mandelbaum, “Weak interactions as manifestations of the substructure of leptons and quarks”, Phys. Lett. B, vol. 102 (5), 319, 1981.
- [33] E. J. Eichten, K. D. Lane, and M. e. Peskin, “New tests for quark and lepton substructure”, Phys. Rev. Lett., vol. 50, 811, 1983.
- [34] N. Cabibbo, L. Maiani, and Y. Srivastava, “Anomalous Z decays: excited leptons?”, Phys. Lett., vol. 139, 459, 1984.
- [35] K. Hagiwara, S. Kmamiya, and D. Zeppenfeld, “Excited Lepton Production at LEP and HERA”, Z. Phys. C, vol. 29, 115, 1985.
- [36] H. Fritzsch, “What I Inside Quarks And Leptons?”, in Proceedings of the Twenty-Second Course of the International School of Subnuclear Physics: Quarks, Leptons and Their Constituents, 1984, Ed. A. Zichichi, New York, Plenum Press, 1988.
- [37] H. Fritzsch, “Composite Quarks And Leptons And Their Flavor Mixing”, in Proceedings of the Europhysics Topical Conference: Flavor Mixing in Weak Interactions, Ed. L.-L. Chau, New York, Plenum Press, 1984.
- [38] G. ‘t Hooft, Recent Developments in Gauge Theories, p. 135, Plenum Press, New York, 1980.
- [39] M. E. de Souza, “The Ultimate Division of Matter”, Scientia Plena, vol. 1(4), 83, 2005.
- [40] M. E. de Souza, “Quark Mass, Quark Compositeness, and Solution to the Proton Spin Puzzle”, Scientia Plena, vol. 4(3), 034801-1, 2008.
- [41] K. Nakamura *et al.*, “Review of Particle Physics: The CKM Quark-Mixing Matrix”, J. Phys. G, vol. 37, 075021, 2010.
- [42] M. E. de Souza, “Weak decays of hadrons reveal compositeness of quarks”, Scientia Plena, vol. 4(6), 064801-1, 2008.
- [43] B. Povh and J. Hüfner, “Systematics of strong interaction radii for hadrons”, J. Phys. Lett. B, vol. 245, 653, 1990.
- [44] T. Koffas, Plenary talk “Status of the Higgs Searches at the LHC: The Channels” at the APS Meeting, April 2013, Denver, CO, USA.
- [45] I. J. R. Aitchison and A. J. G. Hey, “Gauge Theories in Particle Physics”, pp. 259-270, Institute of Physics Publishing, 1993.
- [46] P. D. Morley and S. A. Williams, “On the structure of the alpha particle”, Z. Phys. A, vol. 331(3), 239, 1988.
- [47] M. E. de Souza, “Calculation of almost all energy levels of baryons”, Papers in Physics, vol. 3, 030003, 2011.
- [48] I. V. Gorelov *et al.* (CDF Collaboration), “Evidence for the bottom baryon resonance state Λ_b^{*0} with the CDF II Detector”, hep-ex arXiv:1301.0949v1, Jan 2013, to be published in Proceedings of Science.