

Attraction Magnetic Force Calculation of Axial Passive Magnetic Bearing

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Abstract The paper presents calculation of the force between two ring permanent magnets whose magnetization is axial. Such configuration corresponds to a passive magnetic bearing. Magnets are made of the same material and magnetized uniformly along their axes of symmetry, with the same intensity and same direction. The simple and fast analytical approach is used for this calculation based on magnetization charges and discretization technique. The results for interaction magnetic force obtained using proposed approach are compared with finite element method using FEMM 4.2 software. The proposed algorithm can be used for calculation of microwave devices based on magnetoelectric materials with the use of permanent magnets.

Keywords Permanent Magnet, Interaction Magnetic Force, Magnetization Charges, Discretization Technique, Finite Element Method (FEM)

1. Introduction

Permanent magnets are used nowadays in many applications, and the general need for dimensioning and optimizing leads to the development of calculation methods. Permanent magnets are commonly used in many electrical devices and their own quality depends on the magnet material, magnetization and dimensions. Most engineering applications need several ring permanent magnets and the determination of the magnetic force between them is thus required. Magnetic bearings are contactless suspension devices with various rotating and translational applications[1]. Depending on the ring permanent magnet magnetization direction, the devices work as axial or radial bearings and thus control the position along an axis or the centering of an axis. There are numerous techniques for analyzing permanent magnet devices and different approaches for determining interaction forces between them[2-5]. The authors generally use the Ampere's current model[2,3] or the Columbian approach[4].

2. Theoretical Background

Axial passive magnetic bearing[6] that is considered in the paper is presented in the Figure 1. Since the boundary condition for fictitious surface magnetization charges

density has to be satisfied for both magnets,

$$\eta_{m1} = \hat{n} \cdot \mathbf{M}_1 \text{ and } \eta_{m2} = \hat{n} \cdot \mathbf{M}_2, \quad (1)$$

it is obvious that fictitious surface magnetization charges exist only on the bottom and the top bases of each permanent magnet, because volume magnetization charges for uniform magnetization do not exist. \hat{n} is the unit vector normal to surface.

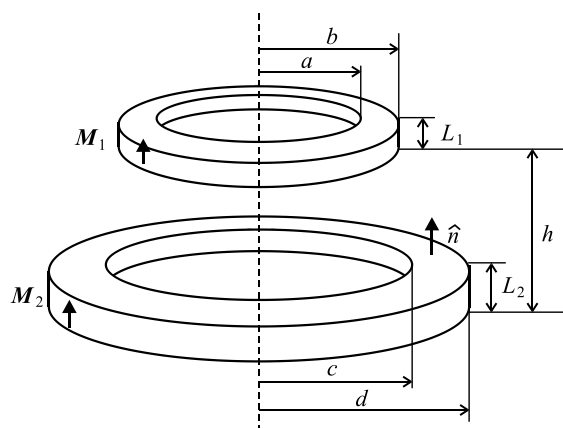


Figure 1. Axial passive magnetic bearing

The simplest procedure for interaction magnetic force determination is to discretize each base of both permanent magnets into system of circular loops[7]. The interaction force between two magnetized circular loops will be calculated first. That will be performed by calculating the magnetic field and magnetic flux density generated by the arbitrary magnetized circular loop of the upper magnet first and then the force that acts on the arbitrary loop of the lower magnet (Figure 2). Using results for interaction magnetic

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force between two circular loops, magnetic force of the axial magnetic bearing can be obtained by summing the contribution of both magnet bases of lower and upper permanent magnets by using uniform discretization technique.

The goal of this approach is to determine the interaction magnetic force between two circular loops uniformly loaded with magnetization charges Q_{m1} and Q_{m2} . Dimensions and positions of the loops are presented in the Figure 2. For determining the interaction force between two circular loops, magnetic scalar potential, magnetic field and magnetic flux density generated by the upper loop will be calculated. Elementary magnetic scalar potential generated by the elementary point magnetization charge, dQ_{m1} , is

$$d\varphi_m = \frac{dQ_{m1}}{4\pi} \frac{1}{R}. \quad (2)$$

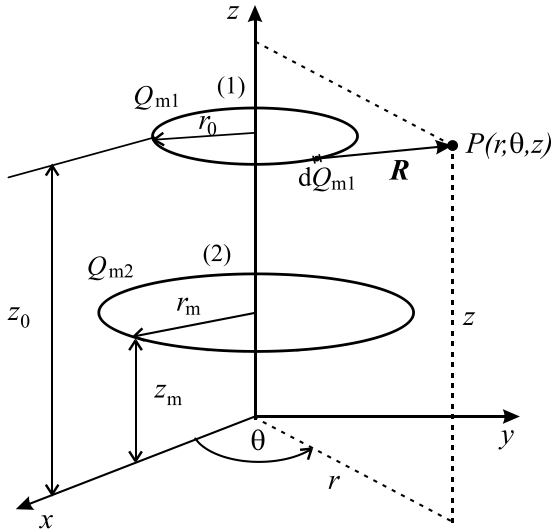


Figure 2. Axial Two circular loops

$$\text{Since } dQ_{m1} = Q'_{m1} dl = \frac{Q_{m1}}{2\pi a} a d\theta' = \frac{Q_{m1}}{2\pi} d\theta',$$

elementary magnetic scalar potential has the following form

$$d\varphi_m = \frac{Q_{m1}}{8\pi^2} \frac{1}{R} d\theta', \quad (3)$$

and the resulting magnetic scalar potential generated by the upper circular loop at an arbitrary point $P(r, \theta, z)$ is

$$\varphi_m = \frac{Q_{m1}}{8\pi^2} \int_0^{2\pi} \frac{1}{\sqrt{r^2 + r_0^2 + (z - z_0)^2 - 2r_0 r \cos(\theta - \theta')}} d\theta', \quad (4)$$

Considering the existing symmetry, in $\theta = 0$ plane, magnetic scalar potential has the following form

$$\varphi_m(r, z) = \frac{Q_{m1}}{4\pi^2} \int_0^\pi \frac{1}{\sqrt{r^2 + r_0^2 + (z - z_0)^2 - 2rr_0 \cos\theta'}} d\theta'. \quad (5)$$

Substituting $\theta' = \pi - 2\alpha$ in Eq. (5), magnetic scalar potential is obtained as:

$$\varphi_m(r, z) = \frac{Q_{m1}}{2\pi^2} \int_0^{\pi/2} \frac{1}{\sqrt{(r + r_0)^2 - 4rr_0 \sin^2(\alpha) + (z - z_0)^2}} d\alpha. \quad (6)$$

After some simple operations the magnetic scalar potential can be given in the form:

$$\varphi_m(r, z) = \frac{Q_{m1}}{2\pi^2} \frac{K\left(\frac{\pi}{2}, k\right)}{\sqrt{(r + r_0)^2 + (z - z_0)^2}} \quad (7)$$

$$\text{where } K\left(\frac{\pi}{2}, k\right) = K = \int_0^{\pi/2} \frac{1}{\sqrt{1 - k^2 \sin^2(\alpha)}} d\alpha,$$

is complete elliptic integral of the first kind with modulus

$$k^2 = \frac{4rr_0}{(r + r_0)^2 + (z - z_0)^2}.$$

External magnetic field (magnetic field generated by the upper loop) at an arbitrary point can be determined as

$$\mathbf{H}^{\text{ext}}(r, z) = -\text{grad } \varphi_m(r, z) = H_r^{\text{ext}}(r, z)\hat{\mathbf{r}} + H_z^{\text{ext}}(r, z)\hat{\mathbf{z}}, \quad (8)$$

External magnetic flux density is

$$\mathbf{B}^{\text{ext}}(r, z) = \mu_0 \mathbf{H}^{\text{ext}}(r, z), \quad (9)$$

$$\mathbf{B}^{\text{ext}}(r, z) = B_r^{\text{ext}}(r, z)\hat{\mathbf{r}} + B_z^{\text{ext}}(r, z)\hat{\mathbf{z}}. \quad (10)$$

with components

$$B_r^{\text{ext}}(r, z) = \mu_0 \frac{Q_{m1}}{2\pi^2} \left(\frac{(r^2 - r_0^2 - (z - z_0)^2) E\left(\frac{\pi}{2}, k\right)}{2r((r - r_0)^2 + (z - z_0)^2) \sqrt{(r + r_0)^2 + (z - z_0)^2}} + \frac{K\left(\frac{\pi}{2}, k\right)}{2r \sqrt{(r + r_0)^2 + (z - z_0)^2}} \right) \quad (11)$$

$$B_z^{\text{ext}}(r, z) = \mu_0 \frac{Q_{m1}}{2\pi^2} \times \frac{(z - z_0) E\left(\frac{\pi}{2}, k\right)}{((r - r_0)^2 + (z - z_0)^2) \sqrt{(r + r_0)^2 + (z - z_0)^2}}, \quad (12)$$

$$\text{where } E\left(\frac{\pi}{2}, k\right) = E = \int_0^{\pi/2} \sqrt{1 - k^2 \sin^2(\alpha)} d\alpha,$$

is complete elliptic integral of the second kind with modulus

$$k^2 = \frac{4rr_0}{(r+r_0)^2 + (z-z_0)^2}.$$

The interaction magnetic force on elementary magnetization charge of lower circular loop

$$dQ_{m2} = Q'_{m2} dl = \frac{Q_{m2}}{2\pi b} b d\beta = \frac{Q_{m2}}{2\pi} d\beta \text{ is } d\mathbf{F} = dQ_{m2} \mathbf{B}^{\text{ext}}(r_m, z_m). \quad (13)$$

Finally, interaction magnetic force components can be expressed as:

$$F_r(r, z) = \mu_0 \frac{Q_{m1} Q_{m2}}{2\pi^2} \left(\frac{(r_m^2 - r_0^2 - (z_m - z_0)^2) E\left(\frac{\pi}{2}, k_0\right)}{2r_m \left((r_m - r_0)^2 + (z_m - z_0)^2\right) \sqrt{(r_m + r_0)^2 + (z_m - z_0)^2}} + \frac{K\left(\frac{\pi}{2}, k_0\right)}{2r_m \sqrt{(r_m + r_0)^2 + (z_m - z_0)^2}} \right) = 0 \quad (14)$$

$$F_z(r, z) = \mu_0 \frac{Q_{m1} Q_{m2}}{2\pi^2} \frac{(z_m - z_0) E\left(\frac{\pi}{2}, k_0\right)}{\left((r_m - r_0)^2 + (z_m - z_0)^2\right) \sqrt{(r_m + r_0)^2 + (z_m - z_0)^2}} = \mu_0 \frac{Q_{m1} Q_{m2}}{2\pi^2} F_{z_p}(r_0, r_m, z_0, z_m) \quad (15)$$

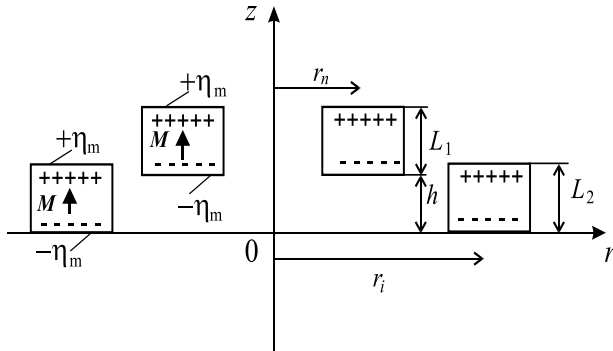


Figure 3. Discretizing model

with elliptic integrals modulus

$$k_0^2 = \frac{4r_0 r_m}{(r_m + r_0)^2 + (z_i - z_m)^2}.$$

The axial component of the force (15) presents interaction force between two magnetized circular loops.

The simplest procedure for attraction magnetic force determination is to discretize each bases of permanent magnets into system of circular loops, where N_1 is the number of discretized segments of each bases of upper permanent magnet and N_2 is the number of discretized segments of each bases of lower permanent magnet.

By taking into account the ring geometry of permanent magnets (Figure 3), the radius of each discretized segment of both bases of upper magnet is

$$r_n = a + \frac{2n-1}{2N_1} (b-a), \quad n=1, 2, \dots, N_1, \quad (16)$$

and magnetization loop charges of upper permanent magnet bases are

$$Q_{mn} = M_1 2\pi r_n \frac{b-a}{N_1}, \quad n=1, 2, \dots, N_1. \quad (17)$$

For lower magnet bases the radius of each discretized segments is

$$r_i = c + \frac{2i-1}{2N_2} (d-c), \quad i=1, 2, \dots, N_2. \quad (18)$$

Magnetization loop charges of lower permanent magnet bases are

$$Q_{mi} = M_2 2\pi r_i \frac{d-c}{N_2}, \quad i=1, 2, \dots, N_2. \quad (19)$$

Using results for interaction magnetic force between two circular loops, Eqs. (15), the attraction magnetic force between two ring permanent magnets can be obtained. It can be achieved by summing the contribution of both magnet bases of lower and upper permanent magnets by using uniform discretization technique,

$$\begin{aligned} F_z = & \frac{2\mu_0 M_1 M_2}{N_1 N_2} (b-a)(d-c) \sum_{n=1}^{N_1} \sum_{i=1}^{N_2} r_n r_i \\ & \times \left(F_{z_p}(r_n, r_i, h, 0) - F_{z_p}(r_n, r_i, h, L_2) \right. \\ & \left. - F_{z_p}(r_n, r_i, h + L_1, 0) + F_{z_p}(r_n, r_i, h + L_1, L_2) \right) \quad (20) \\ F_z = & \frac{2\mu_0 M_1 M_2}{N_1 N_2} (b-a)(d-c) \\ & \times \sum_{n=1}^{N_1} \sum_{i=1}^{N_2} r_n r_i \left(\frac{h E\left(\frac{\pi}{2}, k_1\right)}{\left((r_i - r_n)^2 + h^2\right) \sqrt{(r_i + r_n)^2 + h^2}} + \right. \\ & \frac{(L_2 - h) E\left(\frac{\pi}{2}, k_2\right)}{\left((r_i - r_n)^2 + (L_2 - h)^2\right) \sqrt{(r_i + r_n)^2 + (L_2 - h)^2}} \\ & \left. + \frac{(L_1 + h) E\left(\frac{\pi}{2}, k_3\right)}{\left((r_i - r_n)^2 + (L_1 + h)^2\right) \sqrt{(r_i + r_n)^2 + (L_1 + h)^2}} - \right. \\ & \left. - \frac{(L_1 - h) E\left(\frac{\pi}{2}, k_4\right)}{\left((r_i - r_n)^2 + (L_1 - h)^2\right) \sqrt{(r_i + r_n)^2 + (L_1 - h)^2}} \right) \end{aligned}$$

$$-\frac{(L_2 - L_1 - h)E\left(\frac{\pi}{2}, k_4\right)}{\left((r_i - r_n)^2 + (L_2 - L_1 - h)^2\right)\sqrt{(r_i + r_n)^2 + (L_2 - L_1 - h)^2}} \quad (21)$$

where

$$k_1^2 = \frac{4r_n r_i}{(r_i + r_n)^2 + h^2},$$

$$k_2^2 = \frac{4r_n r_i}{(r_i + r_n)^2 + (L_2 - h)^2},$$

$$k_3^2 = \frac{4r_n r_i}{(r_i + r_n)^2 + (L_1 + h)^2},$$

$$k_4^2 = \frac{4r_n r_i}{(r_i + r_n)^2 + (L_2 - L_1 - h)^2}.$$

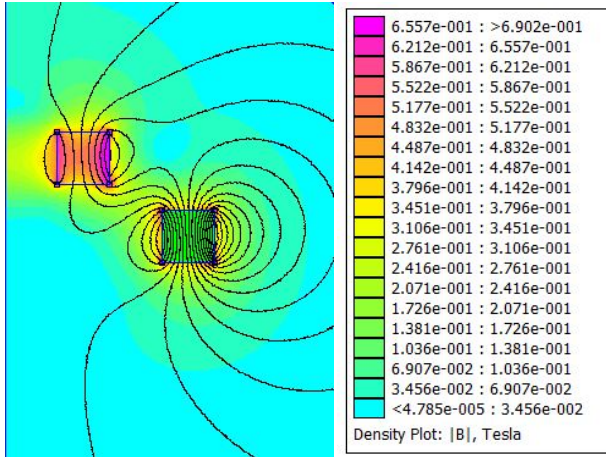


Figure 4. Distribution of magnetic flux density for magnetic bearing (FEMM 4.2 software)

3. Numerical Results

Presumption is that both ring permanent magnets are made of the same material and magnetized uniformly along their axis of symmetry in the same direction, $M_1 = M_2 = M$.

Distribution of magnetic flux density obtained using FEMM 4.2 software[8] is presented in the Figure 4. The values of the geometrical parameters used in the numerical computation are: $a/L_2=1$, $b/L_2=2$, $c/L_2=3$, $d/L_2=4$, $L_1/L_2=1$, $h/L_2=1.5$, $L_2=1$ mm and $M=900$ kA/m.

Convergence of the normalized attraction force,

$$F_z^{\text{nor}} = \frac{F_z}{\mu_0 M^2 L_2^2},$$

obtained using presented approach is given in Table 1 for magnetic bearing dimensions: $a/L_2=1$, $b/L_2=2$, $c/L_2=3$, $d/L_2=4$, $L_1/L_2=1$, $h/L_2=1.5$.

In order to save the calculation time, the number of segments is limited on $N_{\text{tot}} = N_1 + N_2 = 200$ because it is not necessary to take a greater number of segments to obtain a desired accuracy.

Table 1. Convergence of attraction magnetic force versus number of segments

N_{tot}	F_z^{nor}	$F_z^{\text{nor}}(\text{FEM})$
10	-0.304515	-0.306003
20	-0.304502	
30	-0.304497	
50	-0.304494	
100	-0.304493	
200	-0.304492	
300	-0.304492	

Compared results for normalized attraction magnetic force of two identical ring permanent magnets, obtained using presented analytical approach and finite element method (FEM) versus h/L_2 , for parameters: $a/L_2=1$, $b/L_2=2$, $c/L_2=3$, $d/L_2=4$ and $L_1/L_2=1$ are given in the Table 2. Comparative results for normalized interaction magnetic force of axial passive magnetic bearing versus ratios a/L_2 and b/L_2 , obtained using presented approach and finite element method (FEM), for parameters: $c/L_2=3$, $d/L_2=4$, $L_1/L_2=1$ and $h/L_2=1.5$ are shown in the Table 3.

Table 2. Compared results for attraction magnetic force versus h/L_2

h/L_2	F_z^{nor}	$F_z^{\text{nor}}(\text{FEM})$
0	0	0.0000755
0.1	-0.085505	-0.086280
0.2	-0.167360	-0.168159
0.3	-0.242216	-0.243078
0.4	-0.307290	-0.308201
0.5	-0.360536	-0.361493
0.6	-0.400740	-0.401757
0.7	-0.427527	-0.428568
0.8	-0.441304	-0.442424
0.9	-0.443143	-0.444323
1.0	-0.434627	-0.435849
1.1	-0.417663	-0.418920
1.2	-0.394286	-0.395669
1.3	-0.36648	-0.367846
1.4	-0.336043	-0.337446
1.5	-0.304492	-0.306003

Table 3. Compared results for interaction magnetic force versus a/L_2 and b/L_2

a/L_2	b/L_2	F_z^{nor}	$F_z^{\text{nor}}(\text{FEM})$
1.0	2.0	-0.304492	-0.306003
1.5	2.5	-0.432860	-0.434672
2.0	3.0	-0.185029	-0.186999
2.5	3.5	0.950759	0.948785
3.0	4.0	1.89600	1.894140
3.5	4.5	1.105990	1.103435
4.0	5.0	-0.232132	-0.235702
4.5	5.5	-0.657816	-0.662357
5.0	6.0	-0.550606	-0.555822

4. Microwave Devices Based on Magnetolectric Materials with the Use of Permanent Magnets

Magnetolectric (ME) composites that simultaneously exhibit ferroelectricity and ferromagnetism have been of recent research interest due to their potential for applications in multifunctional devices[9].

A main feature of ME devices consists in the application of special permanent magnets[10]. For example, a design of ME microwave phase shifter using the permanent magnets to create of the bias magnetic field is shown in Figure 5.

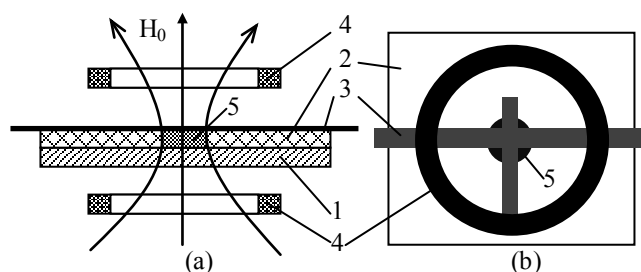


Figure 5. Cross-section (a) and top view (b) ME microwave phase shifter: 1 - basis, 2 - substrate, 3 - microstrip line with a circuit-resonator, 4 - passive magnetic bearings, 5 - ME element

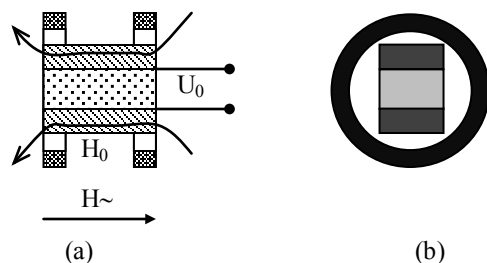


Figure 6. ME sensor of alternating electromagnetic field, longitudinal section (a), lateral view (b)

The ME phase shifter uses the inverse ME effect and operates as follows. Microwave energy passes through the microstrip line with a circuit-resonator. ME element is placed in the center of the substrate under the circuit-resonator. The axial passive magnetic bearing is used for the creation of a permanent bias magnetic field H_0 and for assignment of operating point. Tuning of the device phase is carried out through the set up of the DC voltage between the microstrip line and basis. Smooth change of control voltage allows to smoothly adjust the phase of the output signal in the range from 0 to 360°. Thus, the calculation of the fields created by the passive magnetic bearing necessary for accurate operation of the ME phase shifter.

Another device using the permanent bias magnets is ME low-frequency electromagnetic field sensor[11]. The design of this sensor is shown in Figure 6.

The operation principle of this device is based on the direct ME effect. At placing a sensor in the alternating magnetic field one obtains the alternating electric potential on the the output of the ME element due to direct ME effect. In order to set the correct mode activity of magnetostriction

material in ME element one needs to create a permanent bias magnetic field. This field is created with the help of axial passive magnetic bearing.

5. Conclusions

Determination of the attraction forces of axial passive magnetic bearing is presented. It is preformed using magnetization charges and discretization technique. Presumption was that both magnets are made of the same material and magnetized uniformly along the magnet axis of symmetry, with the same intensity, and same direction. The derived algorithm is easily implemented in any standard computer environment and it enables rapid parametric studies of the attraction force. The results of the presented approach are successfully confirmed using FEMM 4.2 software. Table I shows that it is not necessary to take a great number of segments (not more then 200) to obtain a desired accuracy so the computational time can be saved. Attraction force calculation using presented approach for mentioned parameters and $N_{tot}=200$ is performed with Intel Core 2 Duo CPU at 2.4GHz and 4GB RAM memory and it took less than two seconds of run time. Attraction force is also determined on the same computer using FEMM 4.2 software and the computation time was 14 minutes for about 1.8 million finite elements. Therefore, the advantage of presented analytical approach is its accuracy, simplicity and time efficiency. The proposed algorithm can be used for calculation of the devices based on ME materials with the use of permanent magnets.

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